Description of multipole asymmetries with the Buda Lund hydrodynamical model

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## Introduction

- QGP behaves like perfect fluid $\rightarrow$ hydro description
- Finite number of nucleons $\rightarrow$ generalized geometry is nescessary
- Generalize the space-time and the velocity field distribution
- Higher order flows can be investigated
- HBT radii have $\cos (n \phi)$ dependences in the respective reaction plane
- These can be studied experimentally:

Nucl.Phys. A904-905 (2013) 439c-442c
Phys.Rev.Lett. 112 (2014) 22, 222301


## The Buda-Lund model

## Phys.Rev. C54 (1996) 1390 and Nucl.Phys. A742 (2004) 80-94

- Hydro-model: $S(x, p)=\frac{g}{(2 \pi)^{3}} \frac{p^{\nu} d^{4} \Sigma_{\nu}(x)}{B(x, p)+s_{q}}$ where
$B(x, p)=\exp \left[\frac{p^{\nu} u_{\nu}(x)-\mu(x)}{T(x)}\right]$ is the Boltzmann phase-space distribution and the $p^{\nu} d^{4} \Sigma_{\nu}(x)=p^{\nu} u_{\nu} \delta\left(\tau-\tau_{0}\right) d^{4} x$
- Spatial elliptical asymmetry is ensured by the scaling variable

$$
s=\frac{r_{x}^{2}}{2 X^{2}}+\frac{r_{y}^{2}}{2 Y^{2}}+\frac{r_{z}^{2}}{2 Z^{2}}
$$

- The asymmetry in the velocity field is also elliptical

$$
u_{\mu}=\left(\gamma, r_{x} \frac{\dot{X}}{X}, r_{y} \frac{\dot{Y}}{Y}, r_{z} \frac{\dot{Z}}{Z}\right)
$$

## Generalization of the model I.

- The spatial asymmetry is described by the scaling variable
- General $n$-pole spatial asymmetry (elliptical case: $n=2$ ):

$$
s=\frac{r^{2}}{2 R^{2}}\left(1+\sum_{n} \epsilon_{n} \cos \left(n\left(\phi-\Psi_{n}\right)\right)\right)+\frac{r_{z}^{2}}{2 Z^{2}}
$$

- $\Psi_{n}$ is the angle of the $n$-th order reaction plane



## Generalization of the model II.

- Derive the velocity field from a potential: $u_{\mu}=\gamma\left(1, \partial_{x} \Phi, \partial_{y} \Phi, \partial_{z} \Phi\right)$
- General $n$-pole asymmetrical potential (elliptical case: $n=2$ ):

$$
\Phi=H r^{2}\left(1+\sum_{n} \chi_{n} \cos \left(n\left(\phi-\Psi_{n}\right)\right)\right)+H_{z} r_{z}^{2}
$$

- $u^{\mu} \partial_{\mu} s=0$ is satisfied:
- in $\mathcal{O}\left(\epsilon_{n}\right)$ and $\mathcal{O}\left(\chi_{n}\right)$ if $\dot{\epsilon}_{n}=-2 H \chi_{n}$
- generally there are complicated equations for $\epsilon_{n}(t), \chi_{n}(t)$ functions
- From Euler-equation:
- dynamics of the $\chi_{n} s$ can be calculated
- the initial perturbations are constans during the evolution in first order
- First order calculation can be valid because of the small coefficients


## Observables

- Invariant transverse momentum distribution, flows, azimuthally sensitive HBT radii
- All asymmetries are investigated in their respective reaction plane
- Rotate the system to the second / third order plane and average on the angle of the third / second order plane
- The proper parameters can be set to zero



## Invariant momentum distribution

Significant change could be at high $p_{t}$, the log slope is not affected strongly


## Flows

Elliptic and triangular flows are affected by their own asymmetry parameters


## Mixing of parameters

- The parameters affect the flows together
- The generalization of velocity field is nescessary




## HBT radii

- Calculate in the out - side - long system

$$
R_{\text {out }}^{2}=\left\langle r_{\text {out }}^{2}\right\rangle-\left\langle r_{\text {out }}\right\rangle^{2} \text { and } R_{\text {side }}^{2}=\left\langle r_{\text {side }}^{2}\right\rangle-\left\langle r_{\text {side }}\right\rangle^{2}
$$

where $r_{\text {out }}=r \cos (\phi-\alpha)-\beta_{t} t$ and $r_{\text {side }}=r \sin (\phi-\alpha)$
$\rightarrow$ C. J. Plumberg, C. Shen, U. W. Heinz Phys.Rev. C88 (2013) 044914

- There can be higher order parts
$\rightarrow$ B. Tomášik and U. A. Wiedemann, in QGP3, pp. 715-777.
- We use the following parameterization in
- elliptical case:

$$
R_{\mathrm{out}}^{2}=R_{\mathrm{out}, 0}^{2}+R_{\mathrm{out}, 2}^{2} \cos (2 \alpha)++R_{\mathrm{out}, 4}^{2} \cos (4 \alpha)+R_{\mathrm{out}, 6}^{2} \cos (6 \alpha)
$$

- triangular case:

$$
R_{\mathrm{out}}^{2}=R_{\mathrm{out}, 0}^{2}+R_{\mathrm{out}, 3}^{2} \cos (3 \alpha)+R_{\mathrm{out}, 6}^{2} \cos (6 \alpha)+R_{\mathrm{out}, 9}^{2} \cos (9 \alpha)
$$

- Similar to the $R_{\text {side }}$


## Results of the parametrization - Second order case

This case already have investigated: Eur.Phys.J.A37:111-119,2008 Mainly $\cos (2 \phi)$ behavior but higher order oscillations are also present



## Results of the parametrization - Third order case

Mainly $\cos (3 \phi)$ behavior but higher order oscillations are also present


## Mixing of the parameters

The dependence of the amplitudes of the $R_{\text {out }}^{2}$ and $R_{\text {side }}^{2}$ in the second order case


## Mixing of the parameters

The dependence of the amplitudes of the $R_{\text {out }}^{2}$ and $R_{\text {side }}^{2}$ in the third order case


## Conclusions

- Generalization of $u^{\mu}$ and $s$ with $u^{\mu} \partial_{\mu} s=0$ kept valid
- Higher order flows and azimuthally sensitive HBT radii can be derived
- Absolute value of the azimuthal HBT radii depend on asymmetries
- Third order radii depend on spatial asymmetries more strongly
- Higher order oscillation can be observed in HBT radii
- The spatial and velocity field anisotropies both influence the $v_{n}$ coefficient and the HBT radii
- The asymmetry parameters can be disentangle from the flows and the amplitudes


## Thank you for your attention!

And let me invite you to the 15th Zimanyi School in Budapest http://zimanyischool.kfki.hu/15/

## ZIMÁNYI SCHOOL'15



Arnold Gross: Lexicon
15. Zimányi

## WINTER SCHOOL ON

 HEAVY ION PHYSICSDec. 7. - Dec. 11., Budapest, Hungary


József Zimányi (1931-2006)

## Backup slides - About the spectra

Plot the $N_{1}$ with non zero coefficient divide by $N_{1}$ with zero coefficient


## Backup slides - Higher order amplitudes

## Second order:



## Backup slides - Higher order amplitudes

## Third order:




$$
R_{\text {out, } 9}^{2} / R^{2}{ }_{\text {out }, 0}
$$

$$
R_{\text {side }, 9}^{2} / R_{\text {side }, 0}^{2}
$$




## Backup slides - Averaging

## Averaging vs. set-to-zero



## Backup slides - Square of residuals

An example: square of residuals of $R_{\text {out }}^{2}$ with different parametrizations



Error of $\mathrm{R}_{\text {out, } 2}{ }^{2}$ fit


## Earlier results

Fits with elliptical Buda Lund model: Eur.Phys.J. A47 (2011) 58-66





Value of the parameters

| Meaning | Sign | Value |
| :---: | :---: | :---: |
| Mass of the particle | $m$ | 140 MeV |
| Freeze-out time | $\tau_{0}$ | $7 \mathrm{fm} / \mathrm{c}$ |
| Freeze-out temperature | $T_{0}$ | 170 MeV |
| Temperature-asymmetry parameter | $a^{2}$ | 0.3 |
| Spatial slope parameter | $b$ | -0.1 |
| Transverse size of the source | $R$ | 10 fm |
| Longitudinal size of the source | $Z$ | 15 fm |
| Velocity-space transverse size | $H$ | $10 \mathrm{fm} / \mathrm{c}$ |
| Velocity-space longitudinal size | $H_{z}$ | $16 \mathrm{fm} / \mathrm{c}$ |
| Elliptical spatial asymmetry parameter | $\epsilon_{2}$ | 0.0 |
| Triangular spatial asymmetry parameter | $\epsilon_{3}$ | 0.0 |
| Elliptical velocity-field asymmetry parameter | $\chi_{2}$ | 0.0 |
| Triangular velocity-field asymmetry parameter | $\chi_{3}$ | 0.0 |

Usually one anisotropy parameter is varied, and the others are kept zero

